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DOI:

[10.1002/anie.201700245](https://doi.org/10.1002/anie.201700245)

Document Version

Accepted author manuscript

[Link to publication record in Manchester Research Explorer](#)

Citation for published version (APA):

Kaltsoyannis, N. (2017). Seventeen-coordinate actinide helium complexes. *Angewandte Chemie - International Edition*, 56, 7066–7069. <https://doi.org/10.1002/anie.201700245>

Published in:

Angewandte Chemie - International Edition

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Seventeen-coordinate actinide complexes

Nikolas Kaltsoyannis

Abstract: The geometries and electronic structures of molecular ions featuring He atoms complexed to actinide cations are explored computationally using density functional and coupled cluster theories. A new record coordination number is established, as AcHe_{17}^{3+} , ThHe_{17}^{4+} and PaHe_{17}^{4+} are all found to be true geometric minima, with the He atoms clearly located in the first shell around the actinide. Analysis of AcHe_n^{3+} ($n = 1\text{--}17$) using the Quantum Theory of Atoms in Molecules (QTAIM) confirms these systems as having closed shell, charge-induced dipole bonding. Excellent correlations ($R^2 > 0.95$) are found between QTAIM metrics (bond critical point electron densities and delocalisation indices) and the average Ac–He distances, and also with the incremental He binding energies.

The concept of coordination number (CN) is very important in chemistry. It dates back to 1893, when Alfred Werner defined it as the total number of neighbors around a central atom in a molecule or ion.^[1] Although this definition has been refined – for example to accommodate ligands such as ethene or cyclopentadienyl, which are taken to occupy one and three coordination sites respectively – the concept of CN has stood the test of time.

The current highest CN in molecular chemistry is 16. This has been seen very recently in the experimental work of Pollak *et al.*,^[2] who found 16 Cs...F contacts in $\text{Cs}[\text{H}_2\text{NB}_2(\text{C}_6\text{F}_5)_6]$, and that by Popov *et al.* in CoB_{16}^- ,^[3] and has also been predicted theoretically by Gu *et al.*^[4] for M@Li_{16} ($\text{M} = \text{Ca}, \text{Sr}, \text{Ba}, \text{Ti}, \text{Zr}$ and Hf). The record for the actinide elements stood at 14 for many years (in $\text{U}(\text{BH}_4)_4$),^[5] but is now 15, established experimentally by Daly *et al.* for $\text{Th}(\text{H}_3\text{BNMe}_2\text{BH}_3)_4$,^[6] although the same authors used density functional theory (DFT) to predict that in the gas phase the number of $\text{Th}\cdots\text{H}$ contacts could rise to 16.

All of these contributions make reference to the work of Hermann *et al.*^[7] who, in 2007, published in this journal a DFT survey of the geometric structures of Pb^{2+} coordinated by He atoms. The maximum value of n in PbHe_n^{2+} was found to be 15, with no true minimum structures being located for $n > 15$ in which all the He atoms are located in the first shell around the Pb. Intriguingly, buried at the end of the *Methods* section, Hermann *et al.* write “We note that higher coordination number might be expected for charged actinide–He interactions”. I can now report that this suggestion is correct, and that

computationally stable one-shell AnHe_n^{n+} structures exist for Ac^{3+} , Th^{4+} and Pa^{4+} in which $n = 17$, a new record CN.

As noted by Hermann *et al.*, candidate metals for extensive He coordination should be large and positively charged, and the early actinides in their typical oxidation states certainly fit that description. In order to establish a robust method for exploring PbHe_n^{2+} structures, Hermann *et al.* calculated potential curves for the interaction of Pb^{2+} with a single He atom, and showed that the Generalized Gradient Approximation (GGA) functional PW91, in conjunction with the modest aug-cc-pVDZ basis set, matches very closely the result from high-level *ab initio* coupled cluster (CCSD(T)) calculations using the much larger aug-cc-pVQZ bases.^[7] I have adopted a similar approach, and show in Figure 1 representative results for the parent actinide actinium in its group valence; the binding energy (E_{B}) of Ac^{3+} with He, calculated at both DFT and CCSD(T) levels, is plotted against $r_{\text{Ac-He}}$ from 2 to 5 Å. For a system well described by a single reference electronic configuration, such as He coordinating to a closed shell metal ion, CCSD(T) in conjunction with a large basis set will give essentially exact results. CCSD(T) T_1 diagnostics are very low for AcHe^{3+} (aug-cc-pV5Z for He) at its equilibrium geometry and at 5 Å, 0.003 and 0.002 respectively, confirming the single configuration description.

My results are somewhat different from those for PbHe^{2+} . GGA DFT (PBE) with the aug-cc-pVDZ He basis is now very similar to CCSD(T) with the same small basis set. Increasing the He basis set quality in the CCSD(T) calculations leads, as would be expected, to a deepening of the potential well, giving $E_{\text{B}} = -35.0 \text{ kJ mol}^{-1}$ at 2.552 Å for He with the very large aug-cc-pV5Z basis set. Hence, unlike PbHe^{2+} , GGA DFT with a modest basis set does not well reproduce CCSD(T) with a big basis for AcHe^{3+} . As increasing the He basis quality with PBE makes very little difference to the potential well, I explored many other types of DFT to see if any could replicate the CCSD(T)/aug-cc-pV5Z data. This survey included GGA, hybrid, meta-GGA, meta-hybrid and double hybrid functionals, including where appropriate dispersion corrections, but none matched CCSD(T)/aug-cc-pV5Z well. As is typical, LDA (SVN5) overbinds the system but, amusingly, the venerable exchange-only local spin density Slater functional (here denoted HFS) matches CCSD(T)/aug-cc-pV5Z almost exactly, as shown in Figure 1. I have therefore explored AcHe_n^{3+} structures using the HFS/aug-cc-pVDZ approach and, recognizing that this is an unusual method in modern quantum chemistry, have also conducted analogous calculations with the very widely used PBE functional. As we shall see, the results are similar.

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Supporting information for this article is given via a link at the end of the document.

Figure 1. Potential energy curves for AcHe_n^{3+} using various theoretical methods. In all cases the Ac basis set is the same (the Stuttgart relativistic pseudopotential and accompanying polarised quadruple zeta valence basis set^[9]). The method/He basis is as follows: blue; PBE/aug-cc-pVDZ, grey; CCSD(T)/aug-cc-pVDZ, purple; CCSD(T)/aug-cc-pVQZ, black; CCSD(T)/aug-cc-pV5Z, red; HFS/aug-cc-pVDZ. Analogous CCSD(T)/aug-cc-pV5Z and HFS/aug-cc-pVDZ curves for AcHe_2^{3+} are given in the supporting information.

Table 1. Average, minimum and maximum Ac–He distances (r , Å), and ranges, in HFS AcHe_n^{3+} ($n = 1\text{--}17$).

n	r_{av}	r_{min}	r_{max}	$r_{\text{max}}-r_{\text{min}}$
1	2.564	–	–	–
2	2.570	2.570	2.570	0.000
3	2.576	2.576	2.576	0.000
4	2.575	2.572	2.579	0.007
5	2.573	2.570	2.576	0.007
6	2.588	2.581	2.594	0.014
7	2.588	2.582	2.598	0.016
8	2.592	2.582	2.602	0.020
9	2.593	2.583	2.601	0.018
10	2.599	2.586	2.607	0.021
11	2.603	2.590	2.608	0.017
12	2.607	2.607	2.608	0.002
13	2.617	2.611	2.632	0.021
14	2.626	2.609	2.631	0.023
15	2.637	2.619	2.647	0.027
16	2.654	2.619	2.666	0.047
17	2.671	2.664	2.695	0.031

As with PbHe_n^{2+} , true minimum AcHe_n^{3+} structures can be located for $n = 1\text{--}15$; ball and stick representations of these, and their Cartesian atomic coordinates, are provided for the HFS calculations in the supporting information. As noted above, Hermann *et al.* concluded that if n is increased to 16 or higher, no true minimum structure could be obtained for PbHe_n^{2+} in which all the He atoms are positioned in the first shell around the Pb atom. By contrast, I now report that this is not the case for AcHe_n^{3+} . In this system, both the HFS and the PBE functionals find true minimum first shell structures for $n = 16$ and 17. Ball and stick representations of the HFS structures are shown in Figure 2; to the best of my knowledge, AcHe_{17}^{3+} features the highest CN yet established. As can be seen in Table 1, the average Ac–He distances increase by just over 0.1 Å from $n = 1$ to $n = 17$. This is not unexpected as more He atoms coordinate to the metal center. Even for AcHe_{17}^{3+} , however, the spread of Ac–He distances is very small (0.031 Å), clearly showing that all 17 He atoms are indeed in the first coordination shell. All attempts to coordinate an 18th He atom in the first coordination shell failed.

Figure 2. Optimised (HFS) structures of AcHe_{16}^{3+} (left) and AcHe_{17}^{3+} . Cartesian atomic coordinates are provided in the supporting information.

The incremental He binding energy $|E_{\text{B}}(n)|$ is shown in Figure 2 for both HFS and PBE functionals. It is calculated from equation (1), in which $E(n)$ is the total (SCF) energy of AcHe_n^{3+} and $E(\text{He})$ is the total energy of He. $E_{\text{B}}(n)$ is the energy gain upon attaching a He atom to AcHe_{n-1}^{3+} .

$$E_{\text{B}}(n) = E(n) - E(n-1) - E(\text{He}) \quad (1)$$

The trend for both HFS and PBE is very similar, and is like that reported by Hermann *et al.* for PbHe_n^{2+} . There is a gradual decrease, in absolute terms, to $n = 11$. Hermann *et al.* note that the special stability of the icosahedral $n = 12$ structure leads to a local maximum in the $|E_{\text{B}}(n)|$ graph for PbHe_{12}^{2+} , which is also seen for AcHe_{12}^{3+} at the PBE level. More notable, however, is the reduction in E_{B} at AcHe_{13}^{3+} , particularly for the PBE data, reflecting the disruption of the $n = 12$ complex. Hermann *et al.* attribute the slight increase at $n = 14$ to it being another especially stable structure. After this point the gradient of both lines steepens, as the first coordination shell saturates.

Figure 2. $|E_B(n)|$ for AcHe_n^{3+} ($n = 1-17$). Red: HFS, blue: PBE. The data on which this graph is based, and analogous CCSD(T) data for $n = 1-4$, are given in the supporting information.

For a number of years now, my group has employed the Quantum Theory of Atoms in Molecules (QTAIM)^[9] to study the bonding in 5f compounds,^[10] as well as to quantify the strength of actinide–element bonds.^[11] QTAIM data for AcHe_n^{3+} ($n = 1-17$) are collected in Table 2. Hermann *et al.* note that the Pb^{2+} –He interaction is primarily of the charge-induced dipole type, in which there is only minimal transfer of charge from He to Pb^{2+} . The partial atomic charges in Table 2 indicate that this description is also appropriate for AcHe_n^{3+} ; only very little of the overall 3+ charge is acquired by any given He atom, unsurprising given its extremely high ionisation energy (24.6 eV). Addition of successive He atoms to the complex results in a gradual decrease in the partial charge of the Ac center, and this decrease for successive He addition becomes smaller until a limiting value of c. +2.65 is reached for the largest systems. The very small bond critical point electron densities confirm the closed shell nature of the $\text{Ac}\cdots\text{He}$ interactions, as do the energy densities, which average 0.003 in all seventeen systems. The delocalization indices - QTAIM measures of bond order - are clearly very small. ρ and $\delta_{\text{Ac,He}}$ decrease slightly with n , reflecting the lengthening of the Ac–He distances. Indeed, R^2 for the correlation of the average $r_{\text{Ac-He}}$ (Table 1) with ρ and $\delta_{\text{Ac,He}}$ is 0.996 and 0.953 respectively.

Table 2. QTAIM atomic charges q , average Ac–He bond critical point electron densities ρ (atomic units), and Ac–He delocalisation indices δ , in HFS AcHe_n^{3+} ($n = 1-17$).

n	q_{Ac}	q_{He} (av)	ρ (av)	$\delta_{\text{Ac,He}}$ (av)
1	2.966	0.034	0.021	0.132
2	2.933	0.034	0.020	0.130
3	2.902	0.033	0.020	0.128
4	2.872	0.032	0.020	0.124
5	2.843	0.031	0.020	0.119
6	2.819	0.030	0.019	0.119
7	2.793	0.030	0.019	0.116
8	2.769	0.029	0.019	0.113
9	2.746	0.028	0.019	0.110
10	2.725	0.028	0.019	0.108
11	2.704	0.027	0.019	0.106
12	2.683	0.026	0.018	0.104
13	2.673	0.025	0.018	0.099
14	2.662	0.024	0.018	0.095
15	2.655	0.023	0.017	0.090
16	2.652	0.022	0.017	0.085
17	2.653	0.020	0.016	0.080

The extent to which QTAIM metrics correlate with actinide–element bond strength is very variable, ranging from essentially uncorrelated to highly so. For more detailed information, the reader is directed to references 11. AcHe_n^{3+} provide an opportunity to probe this further in an atypical bonding situation. Correlations of ρ and $\delta_{\text{Ac,He}}$ with $|E_B(n)|$ are very high, with R^2 values of 0.989 and 0.960 respectively. Thus the QTAIM metrics are very good guides to the strength of the closed shell,

charge-induced dipole bonding of He atoms to the spherically symmetric Ac^{3+} . The average Ac–He distance also correlates extremely well with $|E_B(n)|$, $R^2 = 0.989$.

Hermann *et al.* conclude the main text of their contribution by noting “The stability of PbHe_{15}^{2+} suggests that it can be identified by mass spectrometric methods.” The present data indicate a somewhat larger stability for AcHe_{17}^{3+} , but the scarcity and radioactivity of Ac are such that experimental studies are unlikely. I therefore extended my search to the next elements in periodic table, Th, Pa and U, both to explore the generality of the Ac conclusions and also to find more plausible candidates for experimental investigation. For U, no stable one-shell structures were located for UHe_n^{q+} ($q = 2-6$) with $n > 15$. This is also true for Pa^{5+} , but PaHe_{17}^{4+} does retain all of the He atoms in the first coordination shell, albeit with a wider range of metal–He distances than for the Ac analogue (2.47–2.61 Å). The average Pa–He E_B in this system is $-64.0 \text{ kJ mol}^{-1}$, significantly larger than that for AcHe_{17}^{3+} ($-27.1 \text{ kJ mol}^{-1}$) as expected given the smaller charge of the metal ion in the latter. Experimental Pa chemistry is also highly challenging, but that of Th is much less so, and hence I was pleased to find that both ThHe_{17}^{3+} and ThHe_{17}^{4+} are stable structures at the HFS and PBE levels. Cartesian coordinates of AnHe_{17}^{q+} ($\text{An} = \text{Th}$, $q = 3, 4$; $\text{An} = \text{Pa}$, $q = 4$) are provided in the supporting information. ThHe_{17}^{3+} has a much wider range of metal–He distances than either AcHe_{17}^{3+} or PaHe_{17}^{4+} (2.67–2.94 Å), making its description as a genuine CN 17 structure debatable, but the geometry of ThHe_{17}^{4+} is very similar to that of AcHe_{17}^{3+} , with Th–He distances in the range 2.54–2.59 Å, and an average Th–He E_B of $-65.2 \text{ kJ mol}^{-1}$. ThHe_{17}^{4+} is therefore the most likely of the new CN 17 systems to be experimentally accessible.

Computational Details

All DFT calculations were performed with the Gaussian 09 code, revision D.01.^[12] (14s 13p 10d 8f 5g)/[10s 9p 5d 4f 3g] segmented valence basis sets with Stuttgart-Bonn variety relativistic pseudopotentials were used for the actinides,^[8] and the aug-cc-pVXZ ($X = \text{D, Q, 5}$) basis sets were employed for He, as discussed in the main text. For the production calculations, the HFS^[13] and PBE^[14] functionals were used, in conjunction with the ultrafine integration grid. The standard SCF convergence criterion (10^{-8}) was used. All geometry optimisations were performed without symmetry constraints, and the geometry convergence criteria were tightened from the default via IOP(6/7=67), which produces 10^{-4} au for the maximum force. The resulting structures were verified as true minima via harmonic vibrational frequency analysis. QTAIM analyses were performed using the AIMALL program package,^[15] with .wfx files generated in Gaussian 09 used as input. CCSD(T) calculations were performed using the Molpro 2015.1 code, with all default settings.^[16]

Acknowledgements

I am grateful to Dr Hanshi Hu for bringing the work of Hermann *et al.*^[7] to my attention, and to the University of Manchester's Computational Shared Facility for computational resources.

Keywords: coordination number • actinide • density functional • actinium • CCSD(T) • QTAIM

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COMMUNICATION

AcHe_{17}^{3+} , ThHe_{17}^{4+} and PaHe_{17}^{4+} have stable, one-shell true minimum structures with the highest coordination number yet reported.

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